

# Hints of new physics in flavour anomalies



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# Outline

- Basics of flavour physics
- History of flavour anomalies
- Introduction to meson mixing
  - How mixing and anomalies interact
- Introduction to meson lifetimes
  - How lifetimes and anomalies interact
- Future of anomalies

# The What and Why of Flavour Physics

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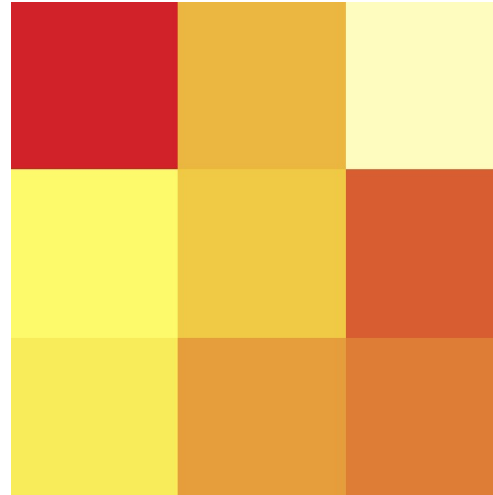
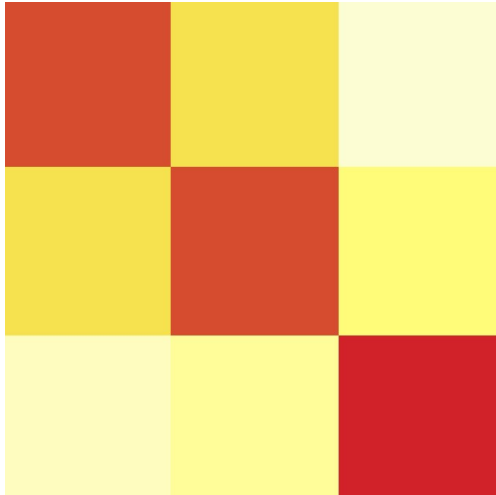
- What is flavour?
  - The different generations of quarks and leptons
- In the SM
  - Only difference is non universal Yukawa coupling to Higgs
    - generates different mass and flavour basis
- Means quarks couple with CKM, leptons with PMNS

# The What and Why of Flavour Physics

- Why? To study these differences – why different masses, why CKM / PMNS look the way they do, why different generations at all?
- Almost easy answer to why 3 generations:
  - Need at least 3 to generate CP violation
- But SM prediction for CP violation off by 10 orders of magnitude from observed baryon asymmetry

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# The What and Why of Flavour Physics

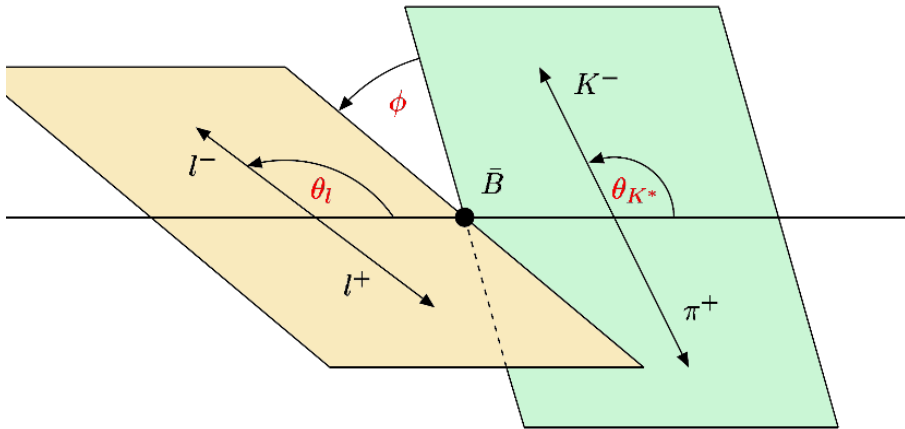
- Why 1: CP violation (big picture)
- Why 2: Lots of flavour changing processes are rare in the SM
  - Easy to enhance, even with high scale NP
- Why 3: Study the SM and our tools
  - Flavour physics is paradigm of EFT – Fermi theory



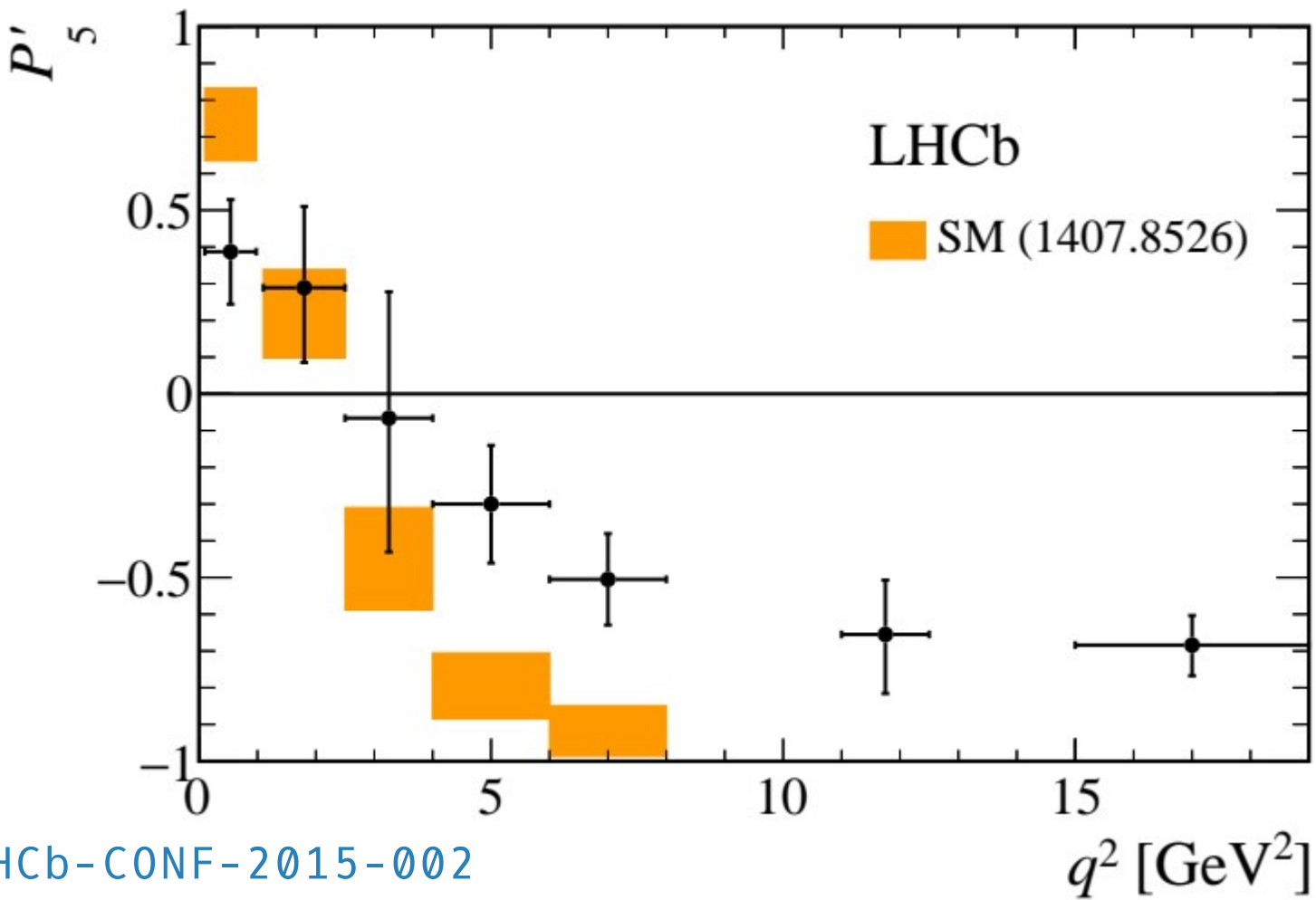
# Flavour anomalies

# Flavour anomalies: a history

- $P'_5$  in 2013,  $3.7 \sigma$  local deviation
- $R_K$  in 2014,  $2.6 \sigma$  local deviation
- $R_{K^*}$  in 2017,  $2-2.5 \sigma$  local deviation

$P_5'$ 


$$\begin{aligned}
 \frac{1}{d\Gamma/dq^2} \frac{d^4\Gamma}{d\cos\theta_\ell d\cos\theta_K d\phi dq^2} = & \frac{9}{32\pi} \left[ \frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell \right. \\
 & - F_L \cos^2 \theta_K \cos 2\theta_\ell + S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi \\
 & + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi \\
 & + S_6 \sin^2 \theta_K \cos \theta_\ell + S_7 \sin 2\theta_K \sin \theta_\ell \sin \phi \\
 & \left. + S_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + S_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right],
 \end{aligned}$$

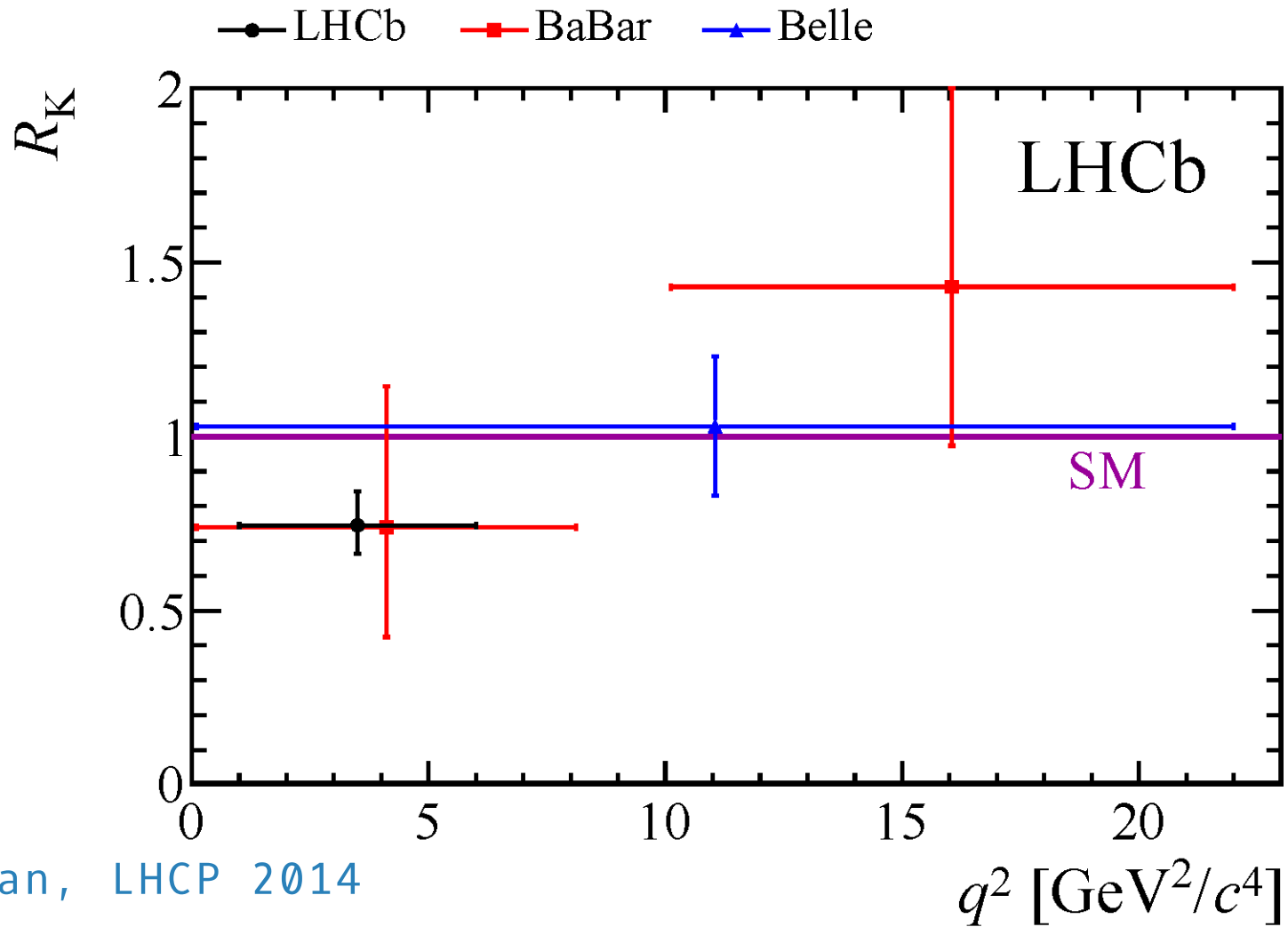


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# Flavour anomalies: a history

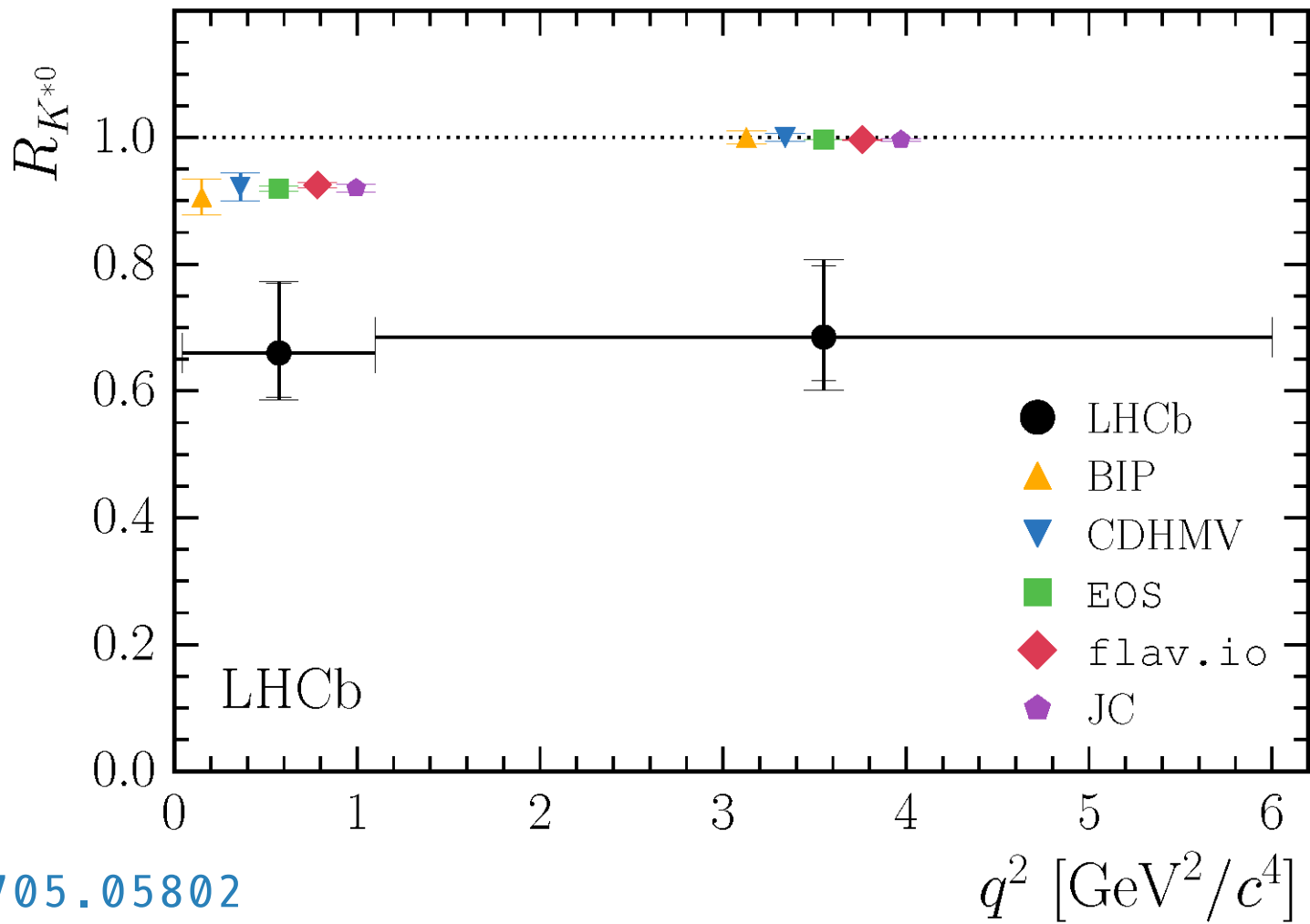
- $P_5'$  in 2013,  $3.7 \sigma$  local deviation
- $R_K$  in 2014,  $2.6 \sigma$  local deviation
- $R_{K^*}$  in 2017,  $2-2.5 \sigma$  local deviation

$$R_{K^{(*)}} = \frac{\mathcal{B}(B \rightarrow K^{(*)} \mu^+ \mu^-)}{\mathcal{B}(B \rightarrow K^{(*)} e^+ e^-)}$$



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$$R_{K^{(*)}}$$

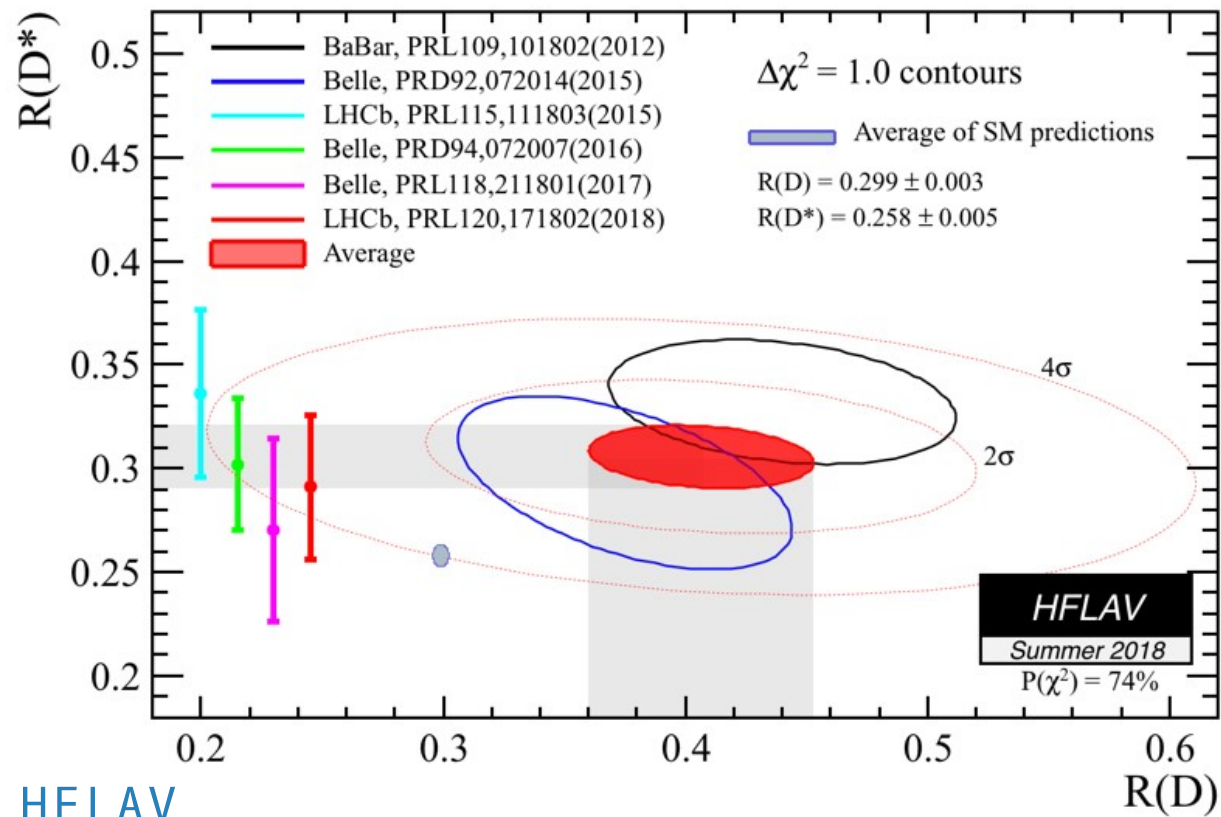
- Very nice as SM predictions are very precise – O(1%)
  - Hadronic uncertainties cancel
    - Note: only in SM – most NP predictions have large uncertainties
- $R_K(1 < q^2 < 6) = 1 \pm 0.01$
- $R_{K^*}(0.045 < q^2 < 1.1) = 0.92 \pm 0.02$
- $R_{K^*}(1.1 < q^2 < 6) = 1 \pm 0.01$

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$$R_{D^{(*)}}$$

- $B \rightarrow D \ell \nu$  decays
- $R_{D^{(*)}} = \text{Br}(B \rightarrow D^{(*)} \tau \nu) / \text{Br}(B \rightarrow D^{(*)} \mu \nu)$
- Tree level, charged current decay
- Overall  $4.1 \sigma$

$$R_D^{(*)}$$



HFLAV

$$R_{D^{(*)}}$$

- $B \rightarrow D \ell \nu$  decays
- $R_{D^{(*)}} = \text{Br}(B \rightarrow D^{(*)} \tau \nu) / \text{Br}(B \rightarrow D^{(*)} \mu \nu)$
- Tree level, charged current decay
- Overall  $4.1 \sigma$
- Not going to talk about this more

)

# Coherent anomalies

- All in  $b \rightarrow s \mu \mu$
- EFT that describes these decays has 6 operators
- Can do global fits to all data, with one or more NP operator in play

$$\begin{aligned} \mathcal{O}_7 &= \frac{e}{16\pi^2} m_b (\bar{s} \sigma_{\mu\nu} P_R b) F^{\mu\nu}, & \mathcal{O}_{7'} &= \frac{e}{16\pi^2} m_b (\bar{s} \sigma_{\mu\nu} P_L b) F^{\mu\nu}, \\ \mathcal{O}_{9\ell} &= \frac{e}{16\pi^2} m_b (\bar{s} \gamma_\mu P_L b) (\bar{\ell} \gamma^\mu \ell), & \mathcal{O}_{9'\ell} &= \frac{e}{16\pi^2} m_b (\bar{s} \gamma_\mu P_R b) (\bar{\ell} \gamma^\mu \ell), \\ \mathcal{O}_{10\ell} &= \frac{e}{16\pi^2} m_b (\bar{s} \gamma_\mu P_L b) (\bar{\ell} \gamma^\mu \gamma_5 \ell), & \mathcal{O}_{10'\ell} &= \frac{e}{16\pi^2} m_b (\bar{s} \gamma_\mu P_R b) (\bar{\ell} \gamma^\mu \gamma_5 \ell). \end{aligned}$$

# Coherent anomalies

- Coherent in the sense that a single NP contribution –  $C_{9\mu}$  – can provide a large improvement in the fit to the data
- With just  $C_{9\mu}$ ,  $5.8 \sigma$  (or  $3.9$  with only LFUV)

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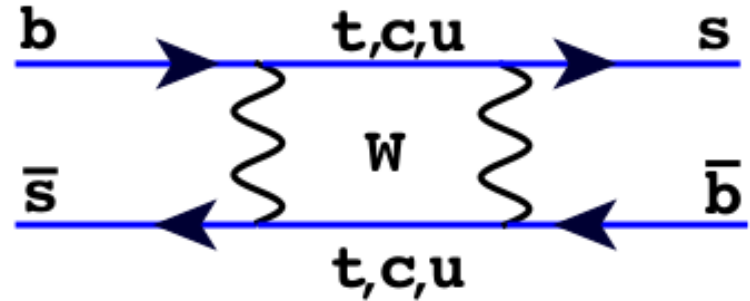
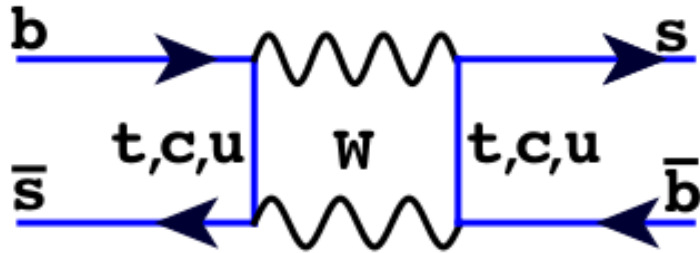


# Meson Mixing

# Short introduction to meson mixing

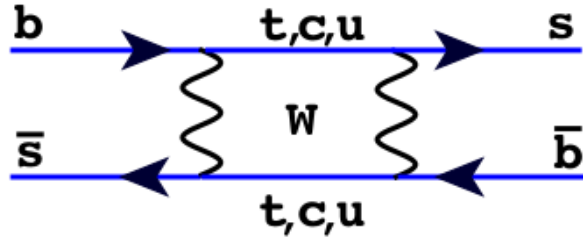
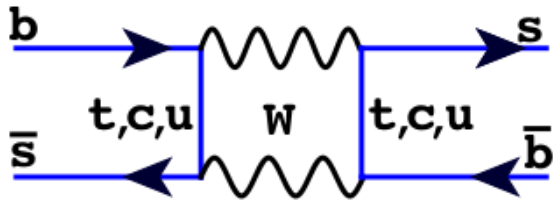
- Consider  $B, \bar{B}$  meson
- Defined by their quark content –  $\bar{b}d, b\bar{d}$ 
  - So they are flavour eigenstates
- But they can oscillate into one another

# Short introduction to meson mixing



# Short introduction to meson mixing

- Can imagine this mixing giving off-diagonal terms in a Schrödinger like equation
- To find mass eigenstates, have to diagonalise



$$\frac{\partial}{\partial t} \begin{pmatrix} B_s \\ \bar{B}_s \end{pmatrix} = \left( \hat{M} - \frac{i}{2} \hat{\Gamma} \right) \begin{pmatrix} B_s \\ \bar{B}_s \end{pmatrix}$$

# Short introduction to meson mixing

- Get two new observables – mass difference and width difference between the two mass eigenstates  $B_H, B_L$  (heavy and light)
- $\Delta M = M_{B_H} - M_{B_L}$
- $\Delta \Gamma = \Gamma_{B_H} - \Gamma_{B_L}$

# Calculating $\Delta M$ and $\Delta \Gamma$

- $\Delta M$  comes from  $\Delta F=2$  operators
- $\Delta \Gamma$  from loop diagrams involving  $\Delta F=1$  operators
  - Because  $\Delta \Gamma$  comes from lifetimes
  - Optical theorem  $\langle B | Q | B \rangle = \text{Im} \sum_X \langle B | Q | X \rangle \langle X | Q | B \rangle$

# Calculating $\Delta M$ and $\Delta \Gamma$

- In the SM, just one operator contributes to  $\Delta M$

$$- (\bar{b}^\alpha \gamma^\mu P_L s^\alpha)(\bar{b}^\beta \gamma_\mu P_L s^\beta),$$

- $\Delta \Gamma$  has many contributing operators

$$(\bar{b}^\alpha \gamma^\mu P_L s^\alpha)(\bar{b}^\beta \gamma_\mu P_L s^\beta),$$

$$(\bar{b}^\alpha P_L s^\alpha)(\bar{b}^\beta P_L s^\beta),$$

$$(\bar{b}^\alpha P_L s^\beta)(\bar{b}^\beta P_L s^\alpha),$$

$$(\bar{b}^\alpha P_L s^\alpha)(\bar{b}^\beta P_R s^\beta),$$

$$(\bar{b}^\alpha P_L s^\beta)(\bar{b}^\beta P_R s^\alpha),$$

# Calculating $\Delta M$ and $\Delta \Gamma$

- $\Delta M \sim C_i \langle Q_i \rangle$ , where  $\langle Q \rangle = \langle B | Q | B \rangle$
- $C_i$  calculated in perturbation theory
- $\langle Q_i \rangle$  need non perturbative technique
  - Lattice QCD
  - Sum rules



$$\langle Q \rangle$$

- Note for later
- For historical reasons,  $\langle Q \rangle$  generally parameterised as  $\langle Q_i \rangle = f_B^2 M_B^2 B_i$
- $B_i$  is bag parameter, contains all the "interesting" physics (assuming you know  $f_B$  already)

Why anomalies  $\rightarrow$  mixing

# Anomalies $\rightarrow$ mixing

- As said earlier, flavour anomalies strongly suggests NP in  $\bar{s}b\bar{\ell}\ell$  operator
- Easy to see that two insertions of NP give  $\bar{s}b\bar{s}b$
- So there is always a link: NP in  $b\rightarrow s\bar{\ell}\ell$  always give NP in  $B_s$  mixing

# Mixing → anomalies

- Reverse is also true
- If we know about mixing, limits what can happen with anomalies
- So what do we know?

# Status of $B_s$ mixing

# $\Delta M_s$ circa 2016

- Experiment:  $17.757 \pm 0.021 \text{ ps}^{-1}$   
HFLAV
- SM:  $18.3 \pm 2.7 \text{ ps}^{-1}$   
1511.09466
  - Relies on FLAG 2013 for  $f_B^2 B$
- SM and experiment in agreement

# $\Delta M_s$ circa 2018

- Experiment:  $17.757 \pm 0.021 \text{ ps}^{-1}$   
HFLAV
- SM:  $20.01 \pm 1.25 \text{ ps}^{-1}$   
1712.06572
  - Relies on FLAG 2017 for  $f_B^2 B$
  - Which is dominated by Fermilab/MILC results from 2016
- SM and experiment disagree at  $\sim 1.8 \sigma$

# $\Delta M_s$ circa 2018

- SM and experiment disagree at  $\sim 1.8 \sigma$
- On its own, not very interesting
- But large class of NP models give positive contribution to  $\Delta M_s$ 
  - i.e.  $\Delta M_s^{\text{th}} \geq \Delta M_s^{\text{SM}}$
  - So  $1.8 \sigma$  discrepancy only gets worse
- (see e.g. [1602.04020](#) for example – CMFV)

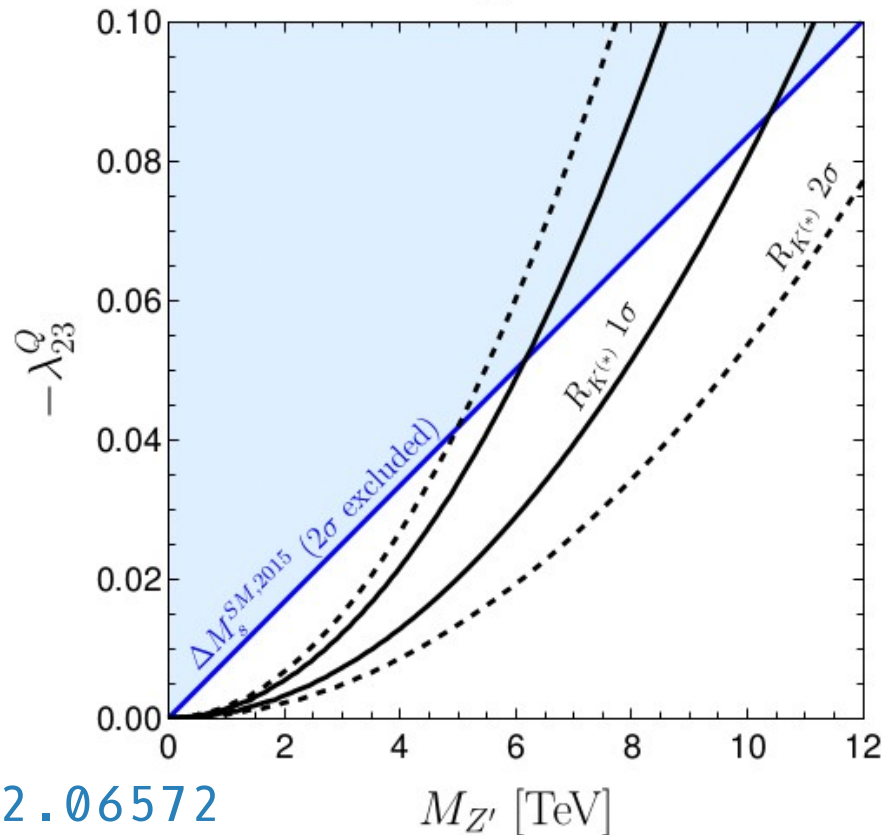


# Concrete example

- Look at how only  $R_{K^{(*)}}$  and  $B_S$  mixing restrict parameter space
- Imagine a new vector boson -  $Z'$
- $Z'_\mu \left( \lambda_{23}^Q \bar{s} \gamma^u P_L b + \lambda_{22}^L \bar{\mu} \gamma^u P_L \mu \right)$

# Concrete example

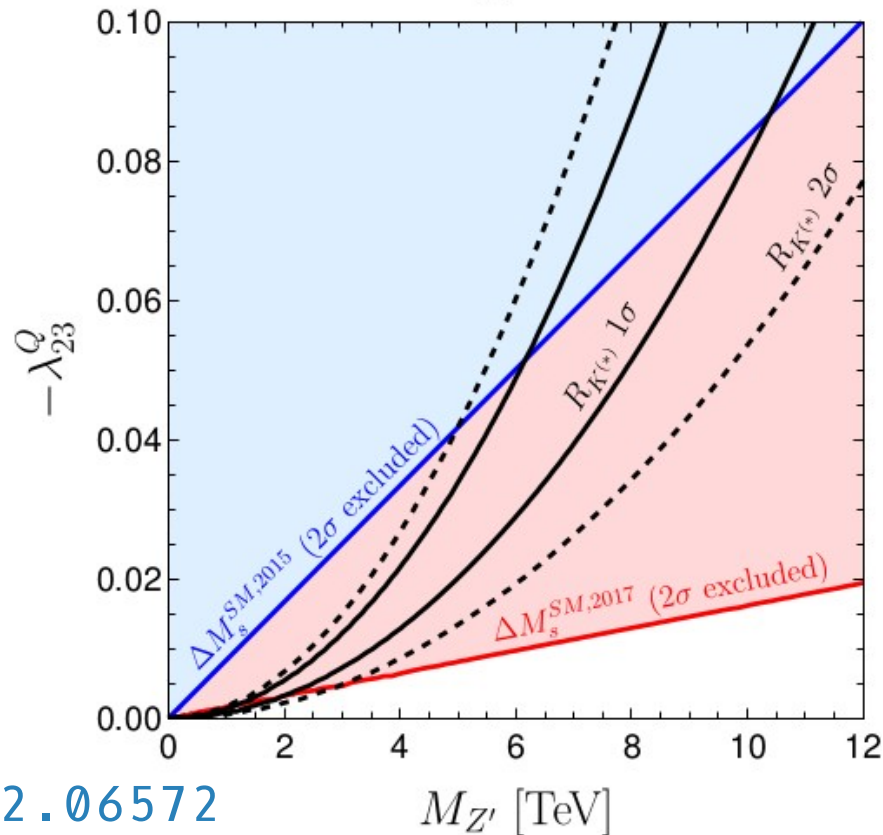
$$\lambda_{22}^L = 1$$



1712.06572

# Concrete example

$$\lambda_{22}^L = 1$$



1712.06572

# Strength of bounds

- Can show that factor of 5 change is generic – applies to any NP model with positive contribution

$$\frac{\Delta M_s^{\text{Exp}}}{\Delta M_s^{\text{SM}}} = \left| 1 + \frac{K}{\Lambda_{\text{NP}}^2} \right| \quad \frac{\Lambda_{\text{NP}}^{2017}}{\Lambda_{\text{NP}}^{2015}} = \sqrt{\frac{\frac{\Delta M_s^{\text{Exp}}}{(\Delta M_s^{\text{SM}} - 2\delta\Delta M_s^{\text{SM}})^{2015}} - 1}{\frac{\Delta M_s^{\text{Exp}}}{(\Delta M_s^{\text{SM}} - 2\delta\Delta M_s^{\text{SM}})^{2017}} - 1}} \simeq 5.2$$

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- Should we believe the new result for  $f_B^2 B$  ?

- Range of different individual numbers

- This is why we average
- In this case, FLAG is the lattice averaging group

Source	$f_{B_s} \sqrt{\hat{B}}$	$\Delta M_s^{\text{SM}}$
HPQCD14	$(247 \pm 12) \text{ MeV}$	$(16.2 \pm 1.7) \text{ ps}^{-1}$
ETMC13	$(262 \pm 10) \text{ MeV}$	$(18.3 \pm 1.5) \text{ ps}^{-1}$
HPQCD09 = FLAG13	$(266 \pm 18) \text{ MeV}$	$(18.9 \pm 2.6) \text{ ps}^{-1}$
<b>FLAG17</b>	$(274 \pm 8) \text{ MeV}$	$(20.01 \pm 1.25) \text{ ps}^{-1}$
Fermilab16	$(274.6 \pm 8.8) \text{ MeV}$	$(20.1 \pm 1.5) \text{ ps}^{-1}$
HQET-SR	$(278^{+28}_{-24}) \text{ MeV}$	$(20.6^{+4.4}_{-3.4}) \text{ ps}^{-1}$
HPQCD06	$(281 \pm 20) \text{ MeV}$	$(21.0 \pm 3.0) \text{ ps}^{-1}$
RBC/UKQCD14	$(290 \pm 20) \text{ MeV}$	$(22.4 \pm 3.4) \text{ ps}^{-1}$
Fermilab11	$(291 \pm 18) \text{ MeV}$	$(22.6 \pm 2.8) \text{ ps}^{-1}$

1712.06572

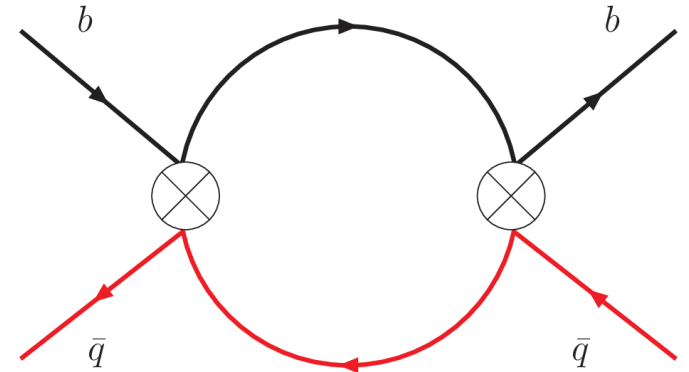
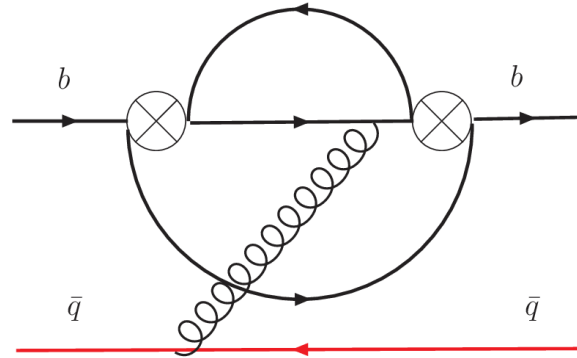
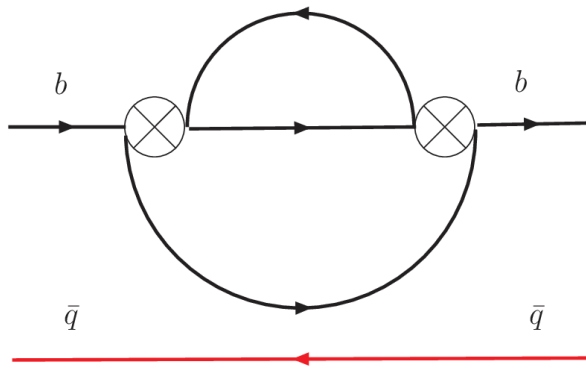
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# Meson lifetimes

# Quick recap on lifetimes

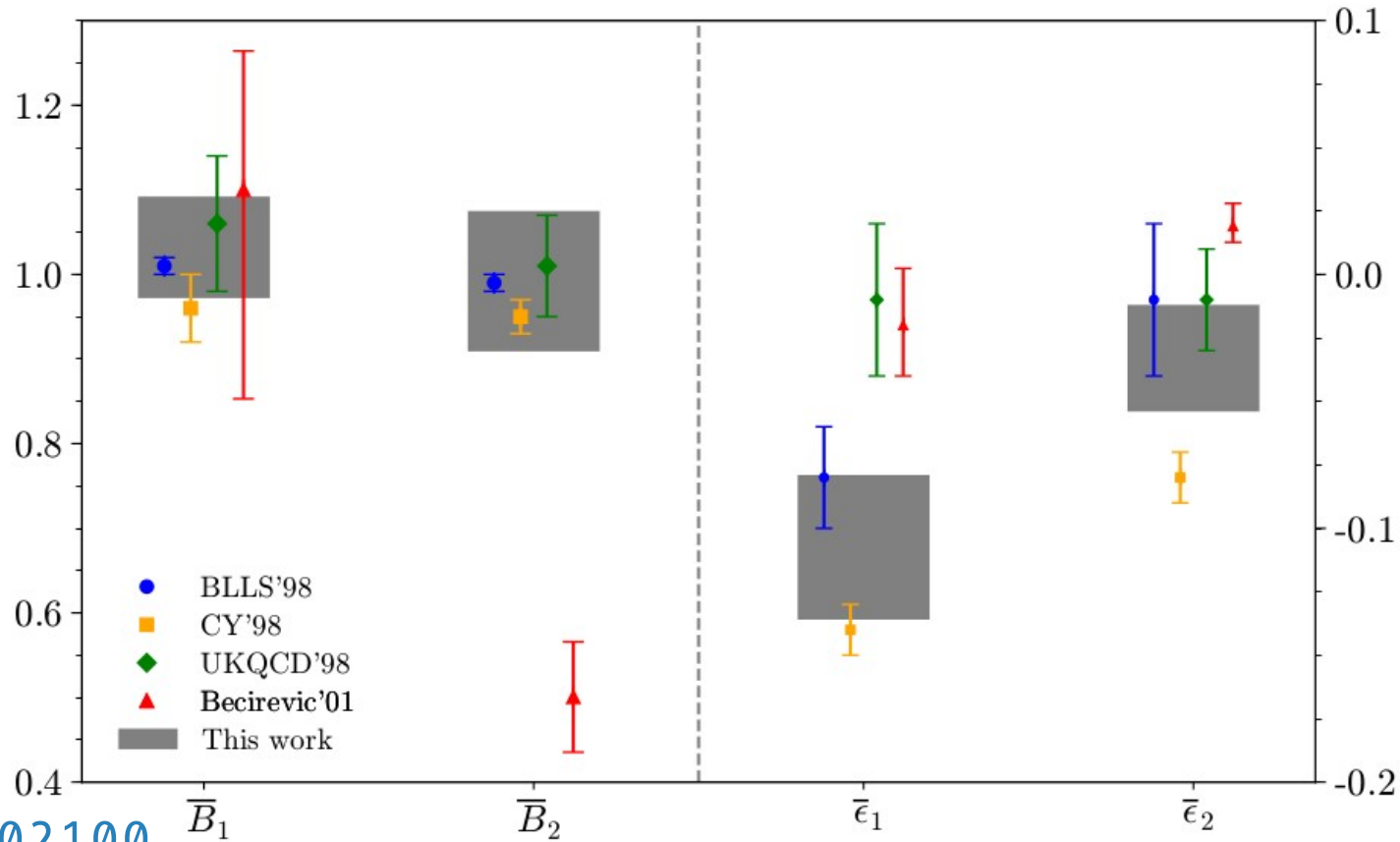
- Use optical theorem to calculate
  - Imaginary parts of  $B \rightarrow B$  processes



# Theory status

- Like mixing, requires hadronic matrix elements to make predictions
- Less well studied by lattice community
- Most recent results from 2001 proceedings
- But recent sum rule calculation also

# Sum rules for bag parameters



1711.02100

# Theory status

- Taking a ratio cancels off various uncertain parameters

- Best theory prediction:  $\frac{\tau(B_s)}{\tau(B_d)} = 1.0005 \pm 0.0011$   
(uncertainty of 0.1%!) 1603.07770

# Lifetime ratio $\tau(B_s)/\tau(B_d)$

- What use is this for the flavour anomalies?
- Most obvious:  $(\bar{s}b)(\bar{\ell}\ell)$  operator contributes to  $B_s \rightarrow \ell\ell$  decay rate  $\rightarrow$  alters lifetime ratio
- However suppressed by  $(m_u/m_b)^2 \simeq 10^{-4}$
- But what about more general NP?

# Lifetime ratio $\tau(B_s)/\tau(B_d)$

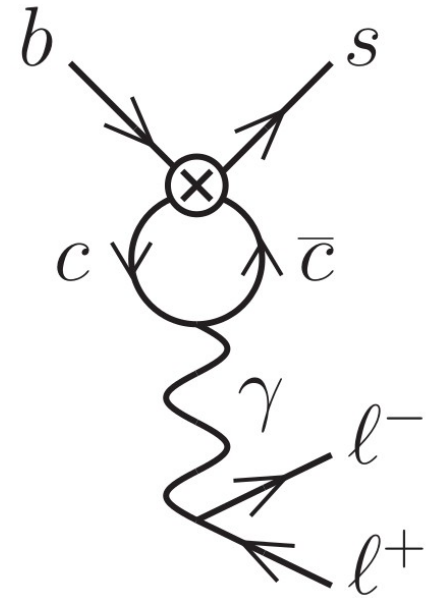
- While LFUV NP is most interesting, seems likely (and fits also support) that there is also contribution that is LFU
  - See e.g. [1704.05446](#), [1809.08447](#)

	Best-fit point	1 $\sigma$ CI	2 $\sigma$ CI
$C_{9\mu}^V$	-1.57	[-2.14, -1.06]	[-2.75, -0.58]
$C_9^U$	0.56	[0.01, 1.15]	[-0.51, 1.78]
$C_{9\mu}^V = -C_{10\mu}^V$	-0.42	[-0.57, -0.27]	[-0.72, -0.15]
$C_9^U$	-0.67	[-0.90, -0.42]	[-1.11, -0.16]

TABLE V. 2D hypotheses. Top: Scenario 7: LFUV and LFU NP in  $C_9^{\text{NP}}$  only. Bottom: Scenario 8:  $C_{9\mu}^V = -C_{10\mu}^V$  and  $C_9^U$  only.

# Lifetime ratio $\tau(B_s)/\tau(B_d)$

- In SM, about half of (LFU) contribution to  $C_9$  comes from charm loops
- So what if NP appears in  $(\bar{s}b)(\bar{c}c)$  ?
- Now lifetime contribution only suppressed by  $(m_c/m_b)^2 \simeq 0.15$

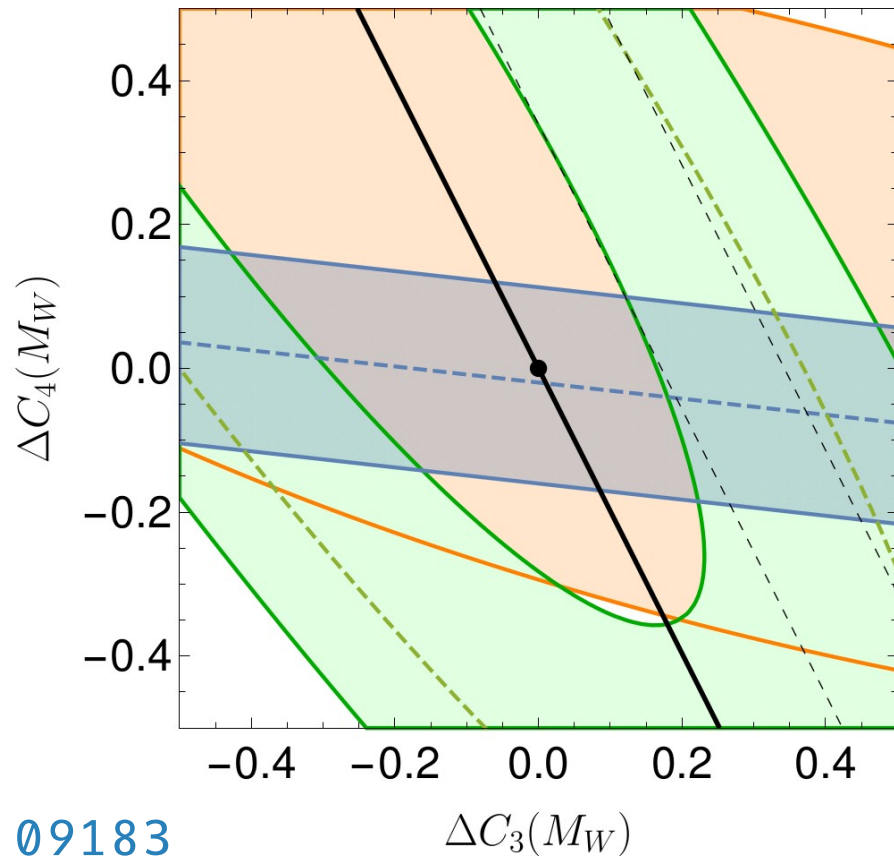
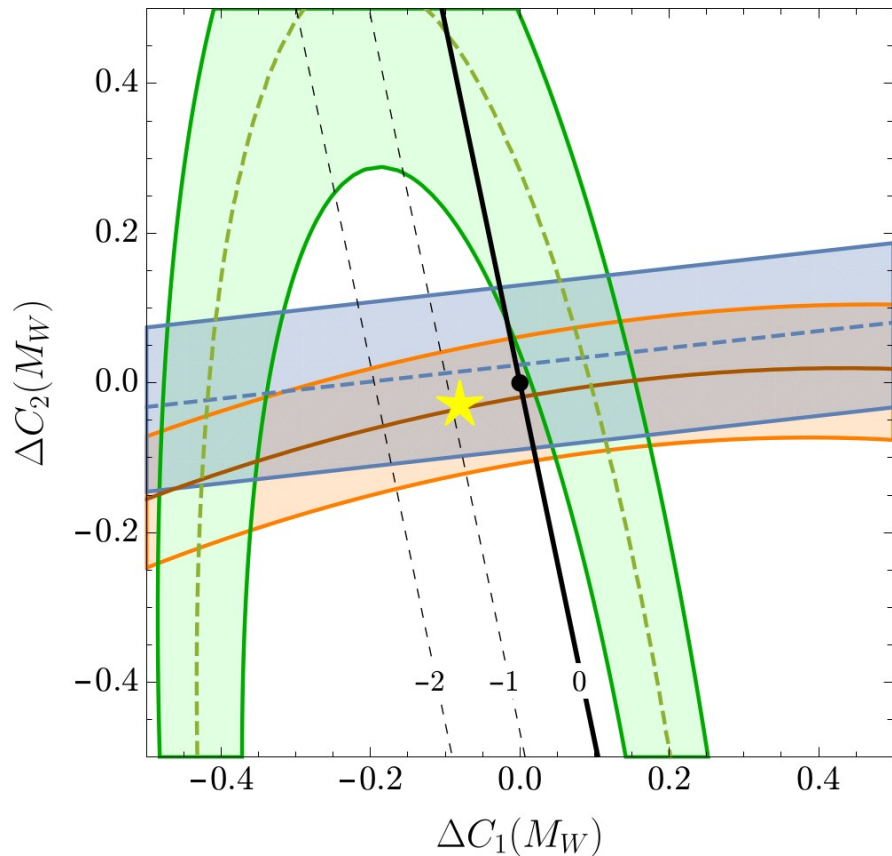




# NP in $(\bar{s} b)(\bar{c} c)$

- Gives rise to correlated effects in several observables
  - Nice way to test, and allows to discriminate between various Dirac structures
- Study in [1701.09183](#) (+ upcoming  $\simeq$  1 month)

# NP in $(\bar{s}b)(\bar{c}c)$



1701.09183

# Future of flavour anomalies

# When will we know?

- Currently, no single measurement has a  $5\sigma$  deviation from SM
  - i.e. no “discovery”
- When might we expect this to happen?
- (Disclaimer – not an experimentalist, numbers taken blindly from their talks)

# LFUV – $R_{K^{(*)}}$

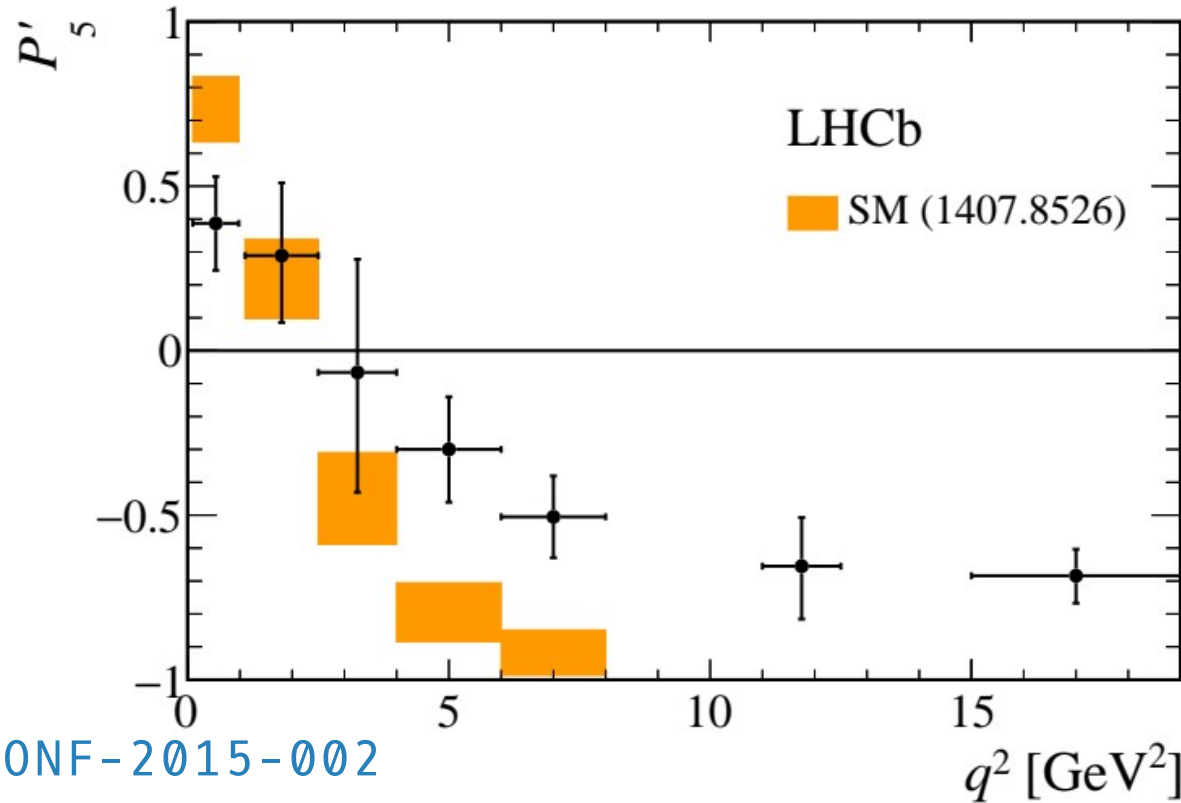
- Now: uncertainty on  $R_{K^{(*)}} \sim 12\%$  (run 1 data)
- In progress, update to  $R_K$  with run 2 data
  - If central value remains the same, 7% uncertainty
- LHCb 2025: Uncertainty 3-4%
  - If same central value  $\rightarrow 10\sigma$  deviation
- Belle II should be able to confirm

# LFUV – $R_K^{(*)}$

Observable	Current LHCb	LHCb 2025	Belle II	Upgrade II
<b>EW Penguins</b>				
$R_K (1 < q^2 < 6 \text{ GeV}^2 c^4)$	0.1 <span style="border: 1px solid green; padding: 2px;">274</span>	0.025	0.036	0.007
$R_{K^*} (1 < q^2 < 6 \text{ GeV}^2 c^4)$	0.1 <span style="border: 1px solid green; padding: 2px;">275</span>	0.031	0.032	0.008
$R_\phi, R_{pK}, R_\pi$	–	0.08, 0.06, 0.18	–	0.02, 0.02, 0.05

LHCb-PUB-2018-009

# Angular observables - $P_5'$



LHCb-CONF-2015-002

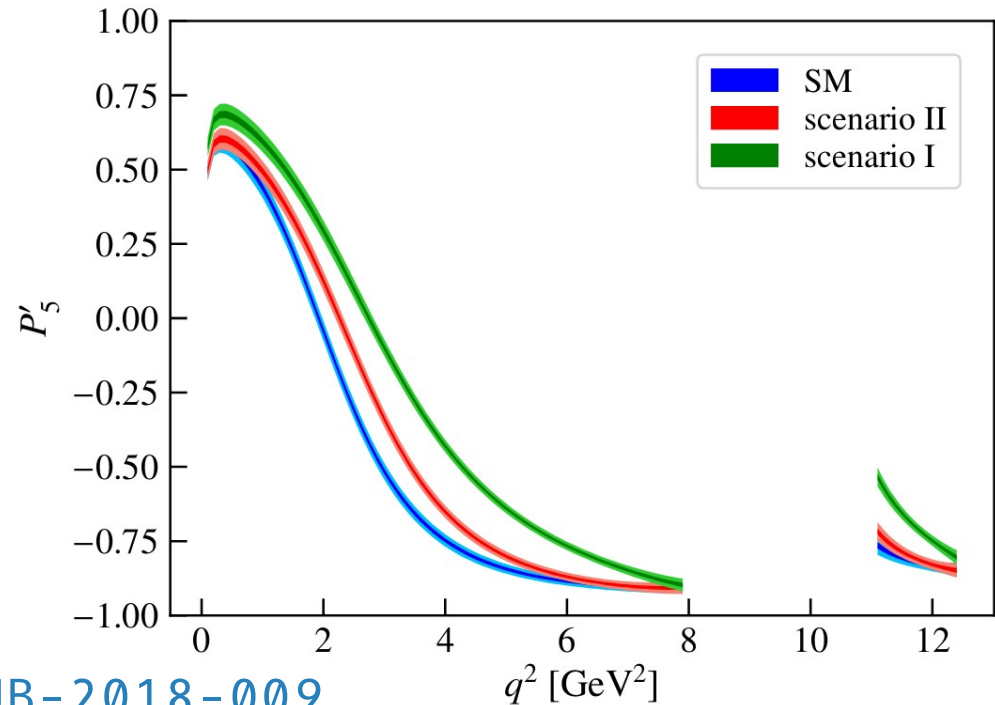
# Angular observables - $P_5'$

- By  $\sim 2035$  (LHCb upgrade 2), can use  $P_5'$  to easily distinguish between various NP scenarios.



# Angular observables - $P_5'$

- Red is  $C_9^\mu = -C_{10}^\mu = -0.7$
- Green is  $C_9^\mu = -1.4$
- Blue is SM
- $3\sigma$  contours

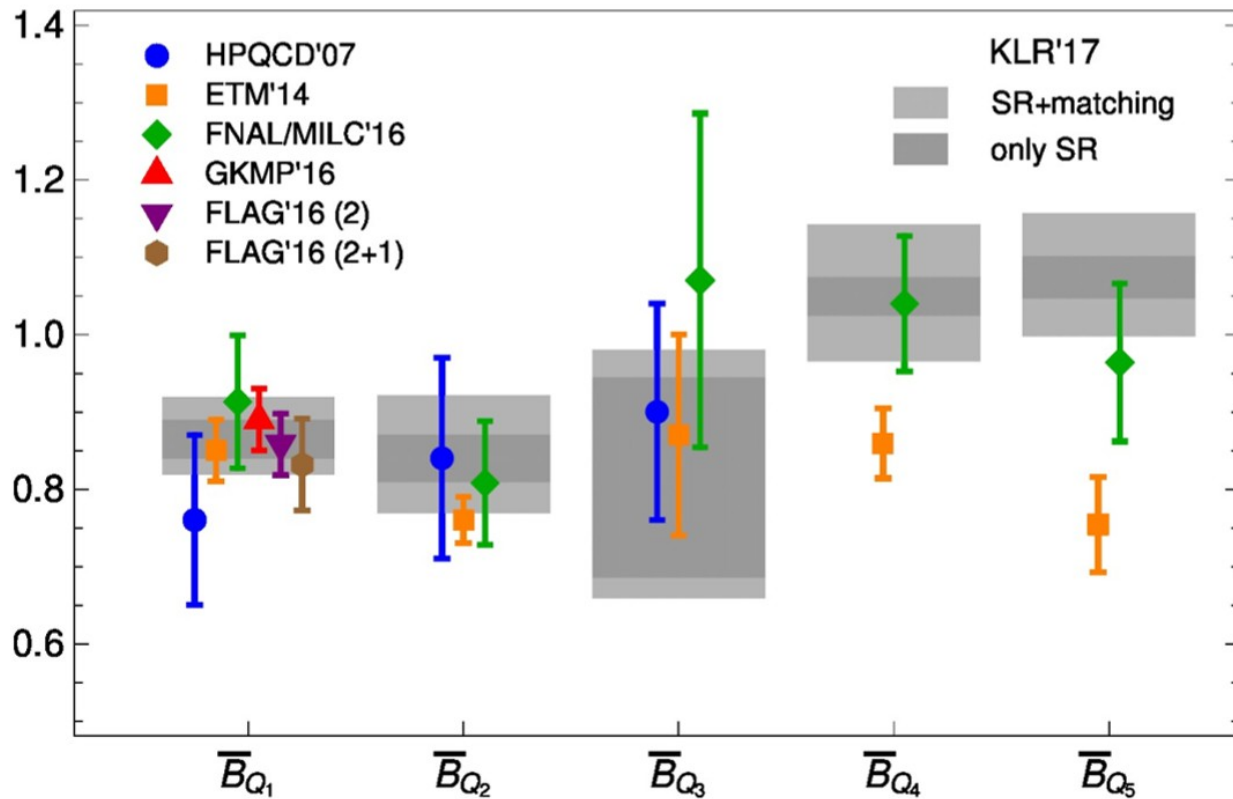


# Summary

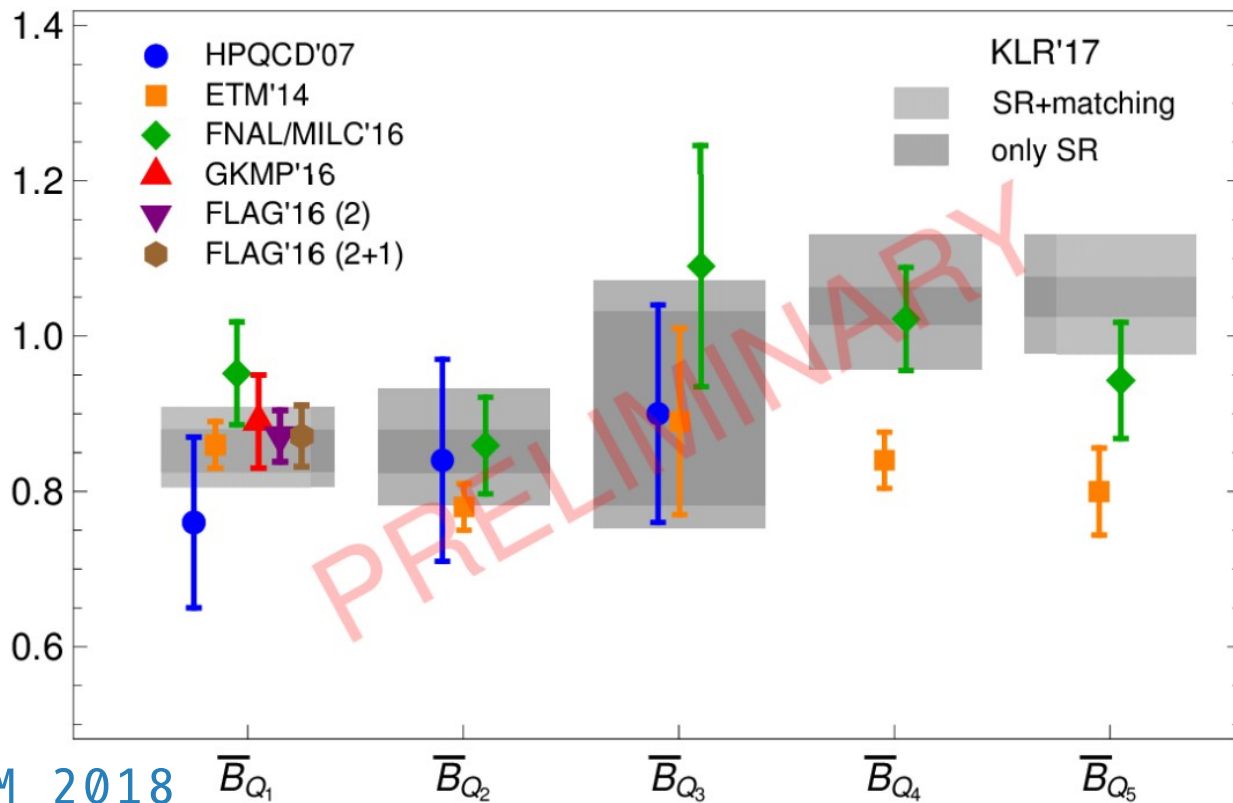
- Flavour anomalies possibly most exciting signs of NP at the moment
  - Unexpected area: LFUV
- Meson mixing very important in constraining BSM models
  - Lattice results the key
- But soon we will know for sure
  - Then a variety of other flavour observables (e.g. lifetimes) will play their part

# Backup

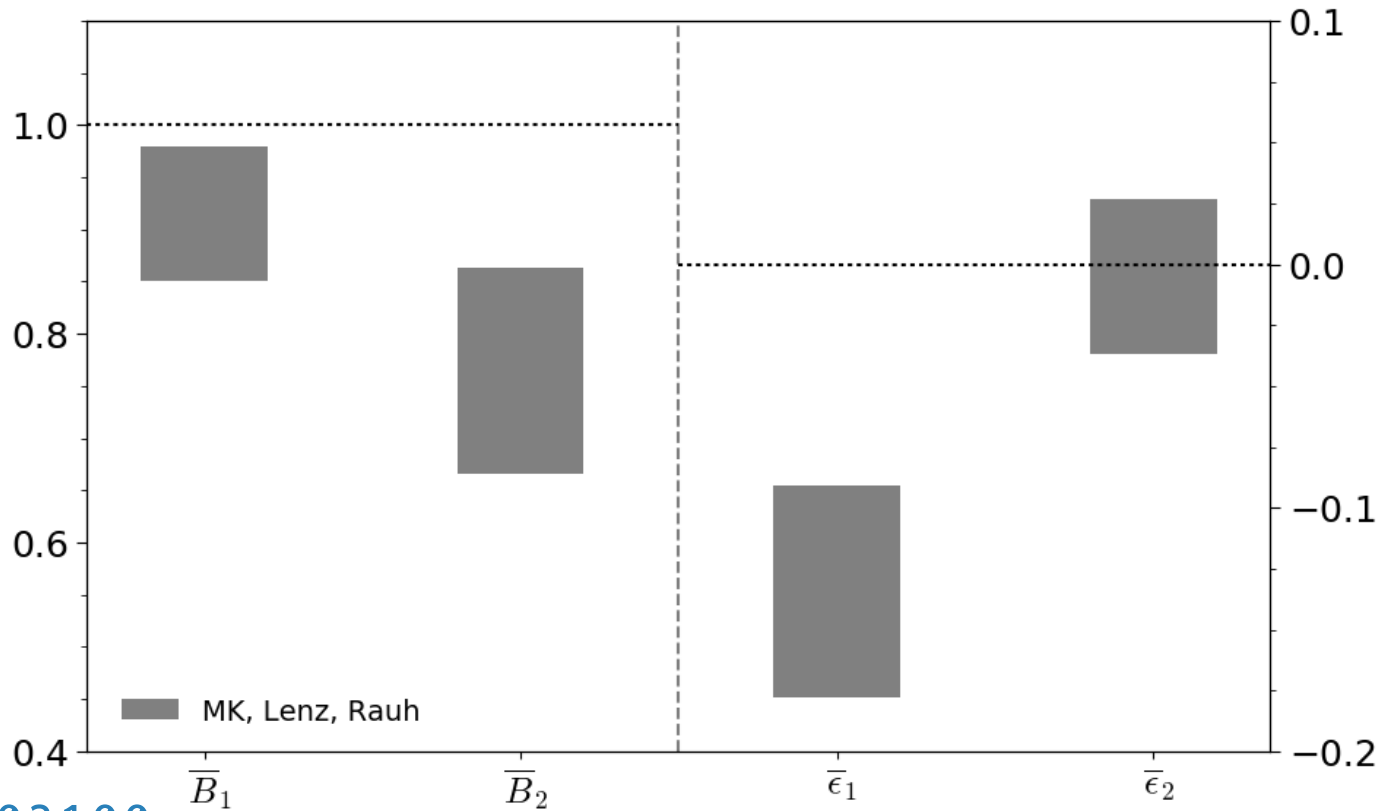
# Sum rules



# Sum rules



# Sum rules



1711.02100

# FLAG discrepancy

- FLAG 2017 average:  $f_{B_s} \sqrt{\hat{B}} = 274 \pm 8 \text{ MeV}$
- But they also give
  - $f_{B_s} = 228.4 \pm 3.7 \text{ MeV}$
  - $\hat{B} = 1.35 \pm 0.06$
- Naive combination:  $f_{B_s} \sqrt{\hat{B}} = 265 \pm 7 \text{ MeV}$

# $V_{cb}$ dependence

