## Vector-like leptons

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(based on 2008.03261 plus 2002.07184/2008.01113 by other authors)

## Vector-like leptons

- Heavy counterpart to SM leptons, but with L \& R having same quantum numbers
- Non-chiral or vector-like


## VLL - motivations

- Can be arbitrarily heavy (not EW scale)
- Unlike new chiral fermions, no strong bounds from Higgs data
- See discussion in 10.1155/2013/910275


## VLL - BSM motivations

- Can arise from composite Higgs, GUTs, neutrino mass models
- E.g. current "hot" topic:
- B anomalies - 4321 model predicts LQ + Z' + colouron + VLQs and VLLs


## Vector-like leptons

- Six possible states (if you want them to couple to the SM)

|  | $S U(3)$ | $S U(2)_{L}$ | $U(1)_{Y}$ |
| :--- | :---: | :---: | :---: |
| $N$ | 1 | 1 | 0 |
| $E$ | 1 | 1 | -1 |
| $\Delta_{1}$ | 1 | 2 | $-1 / 2$ |
| $\Delta_{3}$ | 1 | 2 | $-3 / 2$ |
| $\Sigma_{0}$ | 1 | 3 | 0 |
| $\Sigma_{1}$ | 1 | 3 | -1 |

## LHC bounds

- Direct searches with first generation couplings:
- Singlets: $\mathrm{M}>150 \mathrm{GeV}$
- Doublets: $\mathrm{M}>700 \mathrm{GeV}$
- Triplets: M > 450 GeV

|  | $S U(3)$ | $S U(2)_{L}$ | $U(1)_{Y}$ |
| :--- | :---: | :---: | :---: |
| $N$ | 1 | 1 | 0 |
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## Electroweak precision bounds

- For first generation couplings ( $2 \sigma$ ):

$$
\begin{aligned}
& -N: \mathrm{M} / \mathrm{g}>4 \mathrm{TeV} \\
& -E: \mathrm{M} / \mathrm{g}>5.5 \mathrm{TeV} \\
& -\Delta_{1}: \mathrm{M} / \mathrm{g}>5.5 \mathrm{TeV} \\
& -\Delta_{3}: \mathrm{M} / \mathrm{g}>8 \mathrm{TeV} \\
& -\Sigma_{0}: \mathrm{M} / \mathrm{g}>6 \mathrm{TeV} \\
& -\Sigma_{1}: \mathrm{M} / \mathrm{g}>4 \mathrm{TeV}
\end{aligned}
$$

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& -\Sigma_{1}: \mathrm{M} / \mathrm{g}>2.5 \mathrm{TeV}
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## How do they alter precision physics?

- In general:
- Tree level: modify $Z \ell \ell \& Z \nu \nu$ (Z pole), and $W \ell \nu$ (CKM determination, $G_{F}$ determination)
- 1-loop: modify ( $\bar{\ell} \ell$ ) $(\bar{\ell} \ell)$ contact operator ( $G_{F}$ determination)


## $G_{F}$

- From $\mu \rightarrow e \nu \nu$
- With $(\bar{\ell} \ell)(\bar{\ell} \ell)$ need 1st and 2 nd gen coupling
- With $W \ell \nu$ only need one
- In principle all $V_{u q}$ determinations depend on $G_{F}$ like $V_{u q} \propto 1 / G_{F}$
- But $K_{\mu 2}$ is really $K_{\mu 2} / \pi_{\mu 2}=>G_{F}$ cancels
$G_{F}$



## $G_{F}$



## $G_{F}$

- Cannot reconcile $V_{u s}$ and $V_{u d}$ data with $G_{F}$ alone
- $\delta G_{F} \approx-5 \cdot 10^{-4}$ brings $V_{u d}$ to $V_{u s}^{K_{\mu}}$
- $\delta G_{F} \approx-10^{-2}$ needed to bring $V_{u s}^{K_{\ell 3}}$ to $V_{u s}^{K_{\mu 2}}$
- Factor of 20 difference


## $W \ell \nu$

- Slightly more complex, as $W \ell \nu$ changes directly affect semileptonic decays which determine $V_{u q}$, but also $G_{F}$
$-G_{F} \rightarrow G_{F}\left(1+\delta_{e e}+\delta_{\mu \mu}\right)$
- E.g. NP in $W e \nu$ cancels in beta decay, only sensitive to $W \mu \nu$


## $W \ell \nu$

- $V_{u s} / V_{u d}$ from $K_{\mu 2} / \pi_{\mu 2}$ is independent of both $G_{F}$ and $W \ell \nu$ changes
- $V_{u s}^{K_{\ell 3}}$ sensitive to both $G_{F}$ and $W \ell \nu$, but either only $W e \nu$ or $W \mu \nu$ for $K_{\mu 3}$ or $K_{e 3}$ respectively
- Important to have separate data
$W \ell \nu$
2022

$W \ell \nu$
2022



## Explaining CAA with VLLs



## CKM vs EWPO

- As mentioned, SU2 invariance means changes to $W \ell \nu$ also give changes to $Z \ell \ell$
- So we must test our CKM solutions against EWPO


## CKM vs EWPO



$$
\begin{aligned}
R\left(V_{u s}\right) & \equiv \frac{V_{u s}^{K \mu 2}}{V_{u s}^{\beta}} \equiv \frac{V_{u s}^{K \mu 2}}{\sqrt{1-\left|V_{u d}^{\beta}\right|^{2}-\left|V_{u b}\right|^{2}}} \\
& \approx 1-\left(\frac{V_{u d}}{V_{u s}}\right)^{2} \varepsilon_{\mu \mu} \approx 1-20 \varepsilon_{\mu \mu}
\end{aligned}
$$

## LFV strikes back

- With a single VLL, giving NP in $\mu$ and $e$, you get LFV
- Because $Z \ell_{i} \ell_{j} \sim \sqrt{\left(W \ell_{i} \nu_{i}\right)\left(W \ell_{j} \nu_{j}\right)}$
- And LFV bounds are at least an order of magnitude stronger than other EWPO


## Beyond simplest model

- With two independent VLLs can avoid LFV bounds


## Beyond simplest model

- Consider RH neutrino coupled to electrons, and $\Sigma_{1}$ ( $S U(2)$ triplet equivalent of $\mathrm{RH} e$ ) coupled to muons
- Improves fit by $3 \sigma$



## Conclusions

- VLLs well motivated extensions of the SM
- And can still exist below the TeV scale
- VLLs coupled to muons and electrons can (partially) resolve the CAA
- But EWPO and LFV are important constraints


## Backup

# 4321 VLLs at CMS? 

# Search for pair-produced vector-like leptons in final states with third-generation leptons and at least three $b$ quark jets <br> in proton-proton collisions at $\sqrt{s}=13 \mathrm{TeV}$ 

The CMS Collaboration


#### Abstract

The first search is presented for vector-like leptons (VLLs) in the context of the " 4321 model", an ultraviolet-complete model with the potential to explain existing B physics measurements that are in tension with standard model predictions. The analyzed data, corresponding to an integrated luminosity of $96.5 \mathrm{fb}^{-1}$, were recorded in 2017 and 2018 with the CMS detector at the LHC in proton-proton collisions at $\sqrt{s}=13 \mathrm{TeV}$. Final states with $\geq 3$ b-tagged jets and two third-generation leptons ( $\tau \tau$, $\tau v_{\tau}$, or $v_{\tau} v_{\tau}$ ) are considered. Upper limits are derived on the VLL production cross section in the VLL mass range $500-1050 \mathrm{GeV}$. The maximum likelihood fit prefers the presence of signal at the level of 2.8 standard deviations, for a representative VLL mass point of 600 GeV . As a consequence, the observed upper limits are approximately double the expected limits.


## Beyond simplest model

- My version of Andi's plot



## Beyond simplest model



## LHC bounds

- Direct searches with third generation couplings:
- Singlets: $\mathrm{M}>$ ? GeV
- Doublets: $\mathrm{M}>790 \mathrm{GeV}$
- Triplets: M > ? GeV

|  | $S U(3)$ | $S U(2)_{L}$ | $U(1)_{Y}$ |
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